

# Introduction to physics of quarks with flavour and colour

Alexander Khodjamirian

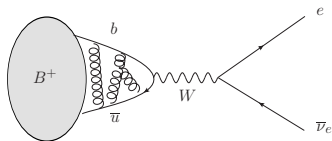


Regional Training Network in Theoretical Physics

Lectures at the Physics Faculty, Yerevan State University, March 30- April 1 2015

## Lecture 3:

### Determination of quark masses and mixing parameters

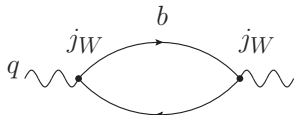


decay amplitude  $\sim V_{ub} \langle 0 | \bar{u} \gamma_\mu \gamma_5 b | B \rangle$   
to calculate **hadronic matrix element**

# Vacuum correlation function

- amplitude of emission and absorption of a  $\bar{u}b$  pair in vacuum by  $j_W = \bar{u}\gamma_5 b$  currents: (using Dirac eq. :  $\partial^\mu(\bar{u}\gamma_\mu\gamma_5 b) = (m_b + m_u)j_W$ )

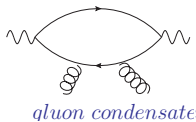
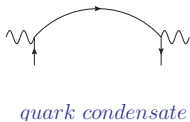
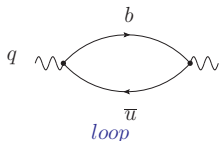
$$\Pi(q^2) = \int d^4x e^{iqx} \langle 0 | T \{ j_W(x) j_W^\dagger(0) \} | 0 \rangle$$



- The 4-momentum  $q = (q_0, \vec{q})$  of the quark-antiquark pair:  
rest frame:  $\vec{q} = 0$ ,  $q^2 = q_0^2$ , fix the energy  $q_0 \ll m_b$
- the  $\bar{u}b$ -pair is **virtual**:  $\Delta E \Delta t \sim 1$ , the energy deficit  $\Delta E \sim m_b$ ,  
 $\Rightarrow$  the quarks propagate during a short time:  $\Delta t \sim 1/m_b$
- the scale  $m_b \gg \Lambda_{QCD}$ : quarks asymptotically free, a loop diagram

# Calculating the correlation function

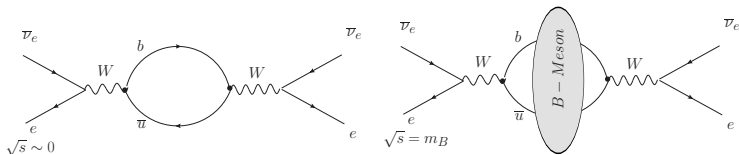
- adding effects of gluon exchange within loop  $\alpha_s(m_B) \ll 1$
- including **nonperturbative** effects due to condensates



- Result: analytical expression for  $\Pi(q^2)$  in terms of  $m_b, m_u$  and **universal** QCD parameters  $\alpha_s, \langle 0 | \bar{q}q | 0 \rangle, \langle 0 | GG | 0 \rangle$

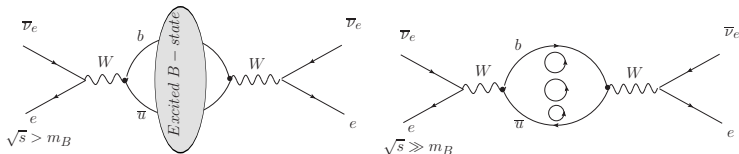
# Relating correlation function with hadrons

- Hypothetical neutrino-electron scattering, varying c.m. energy  $\sqrt{s} = \sqrt{q^2}$ ,
- $\Pi(q^2)$  is the part of the scattering amplitude



highly virtual quark pair,

$B$ -meson, resonance



excited  $B$  mesons

multiple hadrons (continuum)

- $\Pi(q^2)$  can be represented as a **hadronic sum**

# QCD sum rule

- Hadronic sum (dispersion relation)

$$\Pi(q^2) = \frac{\langle 0 | j_W | B \rangle \langle B | j_W | 0 \rangle}{m_B^2 - q^2} + \sum_{B_{exc}, \dots} \frac{\langle 0 | j_W | B_{exc} \rangle \langle B_{exc} | j_W | 0 \rangle}{m_{B_{exc}}^2 - q^2}$$

- derivation based on the unitarity and analyticity of scattering amplitudes in quantum field theory
- using the QCD result for the correlation function:

$$\Pi_{QCD}(q^2; \alpha_s, m_b, m_u, \langle cond \rangle) = \text{Hadronic Sum}$$

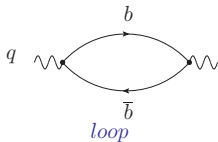
→ QCD sum rule, valid at  $q^2 \ll m_B^2$

- a general method, applicable for various local quark operators and correlation functions
- input:  $\alpha_s(m_Z) = 0.1176 \pm 0.002 \rightarrow \alpha_s(m_B)$   
(perturbative QCD),

# Determination of $b$ -quark mass

- change the current in the correlation function:  $j_W \rightarrow \bar{b}\gamma_\nu b$

no quark condensate,  
only heavy quarks



- the hadronic states are  $\Upsilon$ -meson and excited  $\bar{b}b$  states, well measured in  $e^+e^-$  collisions  $\Rightarrow$  hadronic sum



- fitting  $m_b$  and gluon condensate from the sum rule

$$\Pi^{\bar{b}b}(q^2, \alpha_s, m_b, \langle GG \rangle) = \frac{\langle 0 | \bar{b}\gamma_\mu b | \Upsilon \rangle \langle \Upsilon | \bar{b}\gamma^\mu b | 0 \rangle}{m_\Upsilon^2 - q^2} + \sum_{\Upsilon_{exc}} \frac{\langle 0 | \bar{b}\gamma_\mu b | \Upsilon_{exc} \rangle \langle \Upsilon_{exc} | \bar{b}\gamma^\mu b | 0 \rangle}{m_{\Upsilon_{exc}}^2 - q^2}$$

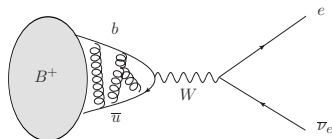
- scale-dependence of the quark mass,  
 $m_b(Q = m_b) = 4.18 \pm 0.03 \text{ GeV}$
- analog  $c$ -quark from charmonium sum rules:  
 $m_c(Q = m_c) = 1.275 \pm 0.025 \text{ GeV}$

# Quark condensate and light quark masses

- correlation functions of light-quark currents  $j_\mu^{ud} = \bar{u}\gamma_\mu(1 - \gamma_5)d$
- input: masses of  $\pi$ -meson, hadronic matrix element  $\langle 0|j^{ud}|\pi\rangle$  from  $\pi \rightarrow \mu\nu_\mu$  decay ( $V_{ud}$  from nuclear  $\beta$  decays)
- the sum rules for strange-quark currents
- results (PDG averages), lattice and non-lattice QCD, with comparable errors:
  - $m_u(Q = 2\text{GeV}) = 2.15 \pm 0.15 \text{ MeV}$ ,
  - $m_d(Q = 2\text{GeV}) = 4.70 \pm 0.20 \text{ MeV}$ ,
  - $m_s(Q = 2\text{GeV}) = 93.5 \pm 2.5 \text{ MeV}$ .
- quark condensate density
  - $\langle 0|\bar{q}q|0\rangle(Q = 2\text{GeV}) = (-277_{-10}^{+12} \text{ MeV})^3$
  - input fixed for the sum rule for  $B$ -meson hadronic matrix element



# Calculation of the $B$ meson decay constant



$$\text{decay constant} \sim \langle 0 | j_W | B \rangle \equiv f_B m_B^2$$

- QCD sum rule (schematically)

$$\Pi_{QCD}(q^2; \alpha_s, m_b, m_u, \langle \bar{q}q \rangle, \dots) = \frac{\langle 0 | j_W | B \rangle \langle B | j_W | 0 \rangle}{m_B^2 - q^2} + \left\{ \sum_{B_{exc}} \right\} \text{duality approx.}$$

- our recent result obtained with NNLO accuracy:  
[P.Gelhausen, D.Rosenthal, AK, A.Pivovarov, Phys.Rev. D88 2013]

$$f_B = 207_{-9}^{+17} \text{ MeV}$$

lattice QCD:  $f_B = 197 \pm 9 \text{ MeV}$  [Fermilab-MILC Collaboration, 2012]

# Extraction of $|V_{ub}|$

- leptonic decay amplitude:

$$A(B^- \rightarrow \tau^- \bar{\nu}_\tau)_{SM} = \frac{G_F}{\sqrt{2}} V_{ub} \langle 0 | \bar{u} \gamma_\mu \gamma_5 b | B \rangle \bar{\tau} \gamma^\mu (1 - \gamma_5) \nu_\tau$$

- partial width: (suppressed for  $\ell = \mu, e$ )

$$BR(B^- \rightarrow \tau^- \bar{\nu}_\tau)_{SM} = \frac{G_F^2 |V_{ub}|^2}{8\pi} m_\tau^2 m_B \left( 1 - \frac{m_\tau^2}{m_B^2} \right)^2 f_B^2 \tau_{B^-},$$

- the most accurate measurement [Belle collaboration (2015)]

$$BR(B \rightarrow \tau \nu_\tau) = (1.25 \pm 0.28(\text{stat.}) \pm 0.27(\text{syst.})) \times 10^{-4}$$

$$\Rightarrow |V_{ub}| = (4.05 \pm 0.39_{\text{exp}} [{}^{+0.18}_{-0.31}] f_B) \times 10^{-3}$$

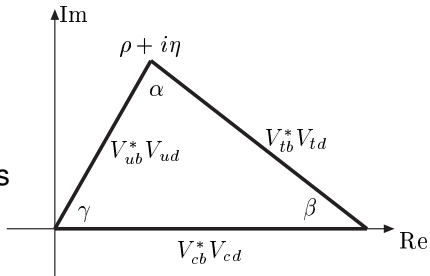
- currently, more accurate  $|V_{ub}|$  from semileptonic  $B \rightarrow \pi \ell \nu_\ell$
- $3\sigma$  discrepancy with "inclusive" determination from  $B \rightarrow X_{u\ell} \nu_\ell$

# $V_{CKM}$ parametrizations and phase

- Unitarity of  $V_{CKM}$ ,  $\lambda \sim \sin \theta_C$ ,

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{1}{2}\lambda^2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4)$$

- $\alpha, \beta, \gamma$   
(over)determined from  
**CP-violating** asymmetries  
in selected nonleptonic  
 $B$ -decays  
(e.g.,  $B \rightarrow J/\psi K \rightarrow \beta$ )



# Summary on $|V_{CKM}|$ determination

$$\left( \begin{array}{ccc} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{array} \right) \Leftrightarrow \left( \begin{array}{ccc} 0^+ \rightarrow 0^+ & K \rightarrow \ell\nu_\ell & B \rightarrow \ell\nu_\ell \\ \text{nuclear } \beta & K \rightarrow \pi\ell\nu_\ell & B \rightarrow \pi\ell\nu_\ell \\ & & B \rightarrow X_u\ell\nu_\ell \\ \\ D \rightarrow \ell\nu_l & D_s \rightarrow \ell\nu_\ell & B \rightarrow D^{(*)}\ell\nu_\ell \\ D \rightarrow \pi\ell\nu_l & D \rightarrow K\ell\nu_\ell & B \rightarrow X_c\ell\nu_\ell \\ \\ B_d \leftrightarrow \bar{B}_d & B_s \leftrightarrow \bar{B}_s & t \rightarrow bW \end{array} \right)$$

- $V_{CKM}$  determination needs **hadronic matrix elements**
- two exceptions:  $|V_{ud}| = 0.97418 \pm 0.00027$   
(superallowed nuclear  $\beta$  transitions, CVC, updated in 2009)

$$|V_{tb}| > 0.74 \quad (\text{single } t\text{-quark decays, CDF+D}\emptyset)$$

$$|V_{tb}| = 0.77^{+0.18}_{-0.24} \quad (\text{precision EW})$$

# Some further reading

available online from <http://inspirehep.net/>

- Y.Grossman , Introduction to flavor physics 1006.3534 [hep-ph] (2010)
- Andrzej J. Buras , Flavour Physics and CP Violation hep-ph/0505175 (2005)
- A. Khodjamirian, ‘Quantum chromodynamics and hadrons: An Elementary introduction,‟ hep-ph/0403145. (2004)